

Modulator Design in String Information Theory

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Abstract

We explore a novel way for information and charge transport in string theory, involving stochastic open-closed transitions of strings interacting with D-branes. Field strings undergoing such transitions act as directional filters, enabling the transport of charged mass strings in a preferred direction. By using a dual filter system and a diffusion mechanism, we demonstrate the emergence of sustained oscillations and charge separation. The resulting dipole field exhibits amplitude modulation, suggesting a string-theoretic realization of radiative signaling via emergent electrodynamic-like behavior. We construct the relevant mathematical model with detailed derivations.

1 Introduction

In conventional string theory, open strings terminate on D-branes while closed strings propagate freely through the bulk. The dynamics of open-closed string transitions have been studied in the context of brane recombination, black hole information, and tachyon condensation. We extend this paradigm by proposing that ambient D-branes can stochastically trigger state transitions of a field string, causing it to alternate between open and closed forms. This behavior is exploited to create directional gating of charged strings, leading to macroscopic charge separation and field generation.

2 Framework

2.1 Open/Closed Transitions

Let a field string $\mathcal{F}$ exist in a D-brane environment such that its state $\sigma(t)$ takes values $\{O, C\}$ for open and closed. The stochastic dynamics of σ are governed by local brane-induced potentials. We define a time-dependent probability $P_O(t)$ that $\mathcal{F}$ is open. This induces a stochastic filter function $f(t)$ defined as

$$f(t) = \begin{cases} 1, & \text{if } \sigma(t) = O, \\ 0, & \text{if } \sigma(t) = C. \end{cases} \quad (1)$$

The transition rates between states are modeled as a Markov process with rates λ_{OC} (closed to open) and λ_{CO} (open to closed), determined by the D-brane tension T_p and string coupling g_s . These rates are approximated as $\lambda_{OC} \sim g_s T_p$ and $\lambda_{CO} \sim g_s^{-1}$, reflecting the energy barriers for string endpoint attachment/detachment on the D-brane.

2.2 Effective Field Derivation

The effective field $\vec{E}_A$ arises from the interaction of an open field string $\mathcal{F}$ with charged mass strings in the presence of a D-brane. Consider a Dp -brane in a 10-dimensional spacetime, with open strings ending on its worldvolume. The interaction between the open string's endpoints and the D-brane gauge field A_μ induces a force on charged mass strings. The effective field is derived from the Born-Infeld action for the D-brane, given by

$$S_{\text{BI}} = -T_p \int d^{p+1}\sigma \sqrt{-\det(g_{\mu\nu} + 2\pi\alpha' F_{\mu\nu})}, \quad (2)$$

where $g_{\mu\nu}$ is the induced metric on the D-brane, $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ is the field strength, and α' is the string scale.

For a field string in the open state ($\sigma(t) = O$), its endpoints couple to A_μ , producing a

localized gauge field. Assuming a simplified configuration where the D-brane lies along the spatial direction x , the effective field acting on a charged mass string with charge q_i is

$$\vec{E}_A = \frac{g_s T_p}{2\pi\alpha'} \hat{x}, \quad (3)$$

where $\hat{x}$ is the unit vector along the D-brane's preferred direction, and the factor $g_s T_p / (2\pi\alpha')$ accounts for the string-brane coupling strength. This field is unidirectional due to the D-brane's geometry, which constrains open string endpoints to move along its worldvolume.

2.3 Charge Filtering and Dipole Formation

In the presence of an open $\mathcal{F}$ -string, charged mass strings experience a directional force, modeled as

$$\vec{F}_i(t) = f(t) q_i \vec{E}_A, \quad (4)$$

where $f(t)$ modulates the force based on the string's state. As $\mathcal{F}$ stochastically switches between open and closed states, charges accumulate asymmetrically along $\pm x$. The position $\vec{x}_i(t)$ of each charged string evolves according to

$$m_i \frac{d^2 \vec{x}_i}{dt^2} = \vec{F}_i(t) - \gamma_{\text{damp}} \frac{d\vec{x}_i}{dt}, \quad (5)$$

where γ_{damp} is a damping coefficient due to interactions with the background.

The net dipole moment is

$$\vec{p}(t) = \sum_i q_i \vec{x}_i(t). \quad (6)$$

In the far-field approximation, the radiative field from the time-varying dipole is

$$\vec{E}_{\text{rad}}(r, t) \sim \frac{1}{r} \ddot{\vec{p}}(t - r/c), \quad (7)$$

where c is the speed of light, and the double derivative accounts for the acceleration of

charges driven by the stochastic force.

3 Type A and B Filters

Two types of filter behaviors are introduced: Type A filters induce transport in the $+x$ direction, while Type B operate in $-x$. These are implemented by orienting D-branes to produce $\vec{E}_A$ in opposite directions. With N such filters (half of each), and a diffusion bridge allowing charge transfer between $+x$ and $-x$ regions, the system forms a self-sustained oscillatory transport cycle. The diffusion bridge is modeled as a region where charges experience a stochastic random walk with diffusion coefficient D , satisfying $\partial_t Q = D \nabla^2 Q$.

4 Stochastic Dynamics and Feedback

The stochastic gating introduces variable transport rates. The charge transport $Q_{\pm}$ in the $\pm x$ regions is governed by

$$\frac{dQ_{\pm}}{dt} = \pm f(t) J(Q_{\mp}) - \gamma Q_{\pm}, \quad (8)$$

where $J(Q) = \kappa Q$ is the current induced by the field, with $\kappa \propto g_s T_p$, and γ is a leakage term due to charge recombination or dissipation. The Markov process for $f(t)$ is solved using the master equation, yielding steady-state probabilities $P_O = \lambda_{OC}/(\lambda_{OC} + \lambda_{CO})$.

5 Amplitude Modulation

The oscillation strength depends on the initial charge Q_0 in the $-x$ region. The dipole moment evolves as

$$\vec{p}(t) \sim A(Q_0) \cos(\omega t), \quad A(Q_0) \propto Q_0, \quad (9)$$

where $\omega \sim \sqrt{\kappa\lambda_{OC}}$ is the oscillation frequency. The emitted radiation is amplitude modulated:

$$\vec{E}_{\text{rad}}(r, t) \sim \frac{-\omega^2 Q_0}{r} \cos(\omega(t - r/c)). \quad (10)$$

6 Discussion

This framework suggests a novel information transmission mechanism within string theory. The open/closed string transition, mediated by D-branes, acts as a dynamic actuator of particle transport. The effective field $\vec{E}_A$, derived from the D-brane's gauge dynamics, connects to established string theory concepts like the Born-Infeld action. By arranging filter pairs and controlling initial conditions, one may engineer emergent radiation with encodable amplitude, potentially relevant to quantum information or gauge/gravity duality.

7 Conclusion

We have proposed a non-traditional mechanism for field generation and information transport in string theory. Using stochastic transitions of field strings and D-brane geometries, charged mass strings are directionally filtered to produce time-varying dipoles. The emergent radiated fields admit amplitude modulation, hinting at communication or computation at a string-theoretic level.

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